Normalized solutions for a quasilinear Schrödinger Choquard equation with exponential critical growth in \mathbb{R}^2

Abstracts: In this paper, we are concerned with normalized solutions to the following quasilinear Schrödinger Choquard equation

$$-\Delta u - u\Delta u^2 + \lambda u = (I_{\alpha} * F(u)) f(u), \text{ in } \mathbb{R}^2,$$

with prescribed mass

$$\int_{\mathbb{D}^2} |u|^2 \ dx = a^2,$$

where a > 0, $\lambda \in \mathbb{R}$, $\alpha \in (0,2)$, I_{α} denotes the Riesz potential, * denotes the convolution opertor, and the nonlinearity f has an exponential critical growth in the sense of Trudinger-Moser inequality. Using Perturbation method and variational methods with Pohozaev manifold, we can avoid the nondifferentiability of the quasilinear term $u\Delta u^2$ and prove the existence of normalized solutions with some further assumption.

Keywords: Normalized solutions; Quasilinear Schrödinger equation; Choquard equation; Exponential critical growth.

1 Introduction

The following generic quasilinear problems have described several physical situations of the form

$$\begin{cases}
i\partial_t \psi = -\Delta \psi - \psi l'(|\psi|^2) \Delta l(|\psi|^2) + f(|\psi|^2) \psi, & \text{in } \mathbb{R}^+ \times \mathbb{R}^2, \\
\psi(0, x) = \psi_0(x), & \text{in } \mathbb{R}^2,
\end{cases}$$
(1.1)

where l, f are given functions, i denotes the imaginary unit, and $\psi: \mathbb{R}^+ \times \mathbb{R}^2 \to \mathbb{C}$ is a complex real function. It is well known that when $l(s) = \sqrt{s}$, problem (1.1) appears in plasma physics and fluid mechanics [17, 27], also in the theory of Heisenberg ferromagnet and in condensed matter theory [23], and the dynamic properties are closely linked to l and f. While, in this article, we focus on the particular case l(s) = s, that is

$$\begin{cases} i\partial_t \psi = -\Delta \psi - \psi \Delta(|\psi|^2) + f(|\psi|^2)\psi, & \text{in } \mathbb{R}^+ \times \mathbb{R}^2, \\ \psi(0, x) = \psi_0(x), & \text{in } \mathbb{R}^2. \end{cases}$$
(1.2)

A stationary wave solution is a solution of the form $\psi(t,x) = e^{-i\lambda t}u(x)$, where $\lambda \in \mathbb{R}$ is a parameter and $u(x) : \mathbb{R}^2 \to \mathbb{R}$ is a time-independent function to be founded. Substitute $\psi(t,x) = e^{-i\lambda t}u(x)$ into (1.2), we obtain the following stationary equation

$$-\Delta u - u\Delta u^2 = \lambda u + f(u), \quad \text{in} \quad \mathbb{R}^2.$$
 (1.3)

For some fixed values of λ , a nontrivial solution of (1.3) is obtained as a critical point of the functional $J_{\lambda}: H^{1,2}(\mathbb{R}^2) \to \mathbb{R}$, which is given by

$$J_{\lambda}(u) := \frac{1}{2} \int_{\mathbb{R}^2} |\nabla u|^2 \, dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 \, dx - \frac{1}{2} \int_{\mathbb{R}^2} \lambda |u|^2 dx - \int_{\mathbb{R}^2} F(u) \, dx,$$

where $F(t) = \int_0^t f(s) ds$, the primitive function of f(t), on the natural space

$$H = \left\{ u \in H^{1,2}(\mathbb{R}^2) : \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 \, dx < +\infty \right\},\,$$

another important way to find the nontrivial solutions for (1.3) is to search for solutions with prescribed mass, that is

$$\begin{cases}
-\Delta u - u\Delta u^2 = \lambda u + f(u), & \text{in } \mathbb{R}^2, \\
\int_{\mathbb{R}^2} |u|^2 dx = a^2,
\end{cases}$$
(1.4)

in this case, $\lambda \in \mathbb{R}$ is part of the unknown. Moreover, it can be obtained by looking for critical points of the corresponding energy functional:

$$I(u) = \frac{1}{2} \int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx - \int_{\mathbb{R}^2} F(u) dx$$
 (1.5)

on the L^2 -sphere

$$\widetilde{S}(a) = \left\{ u \in H : \int_{\mathbb{R}^2} |u|^2 dx = a^2 \right\},$$

which has particular difficulties. To derive the Palais-Smale sequence, one needs new variational methods, because the derived Palais-Smale sequence may not be bounded; even if the Palais-Smale sequence is bounded, the weak limit may not be contained in the L^2 -sphere (even in the radical case). Such difficulties make the study of normalized solutions for (1.4) much more complicated than (1.3) with prescribed $\lambda \in \mathbb{R}$.

One quasilinear term

$$V(u) := \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 \ dx$$

in (1.5) has put forward a new problem: it is not differentiable in the space H. To overcome this difficulty, during the last ten years, various arguments have been put forward on standing wave solutions, while very few results are known about equations of the normalized solutions. Using the minimization methods [26], Nehari manifold approach [18], change variables [9,20] methods, and perturbation method in a series of paper [19,21,22], that recovers the differentiability by considering a perturbed functional

on a smaller function space, one can obtain the standing wave soulutions but not normalized solution. To the best of our knowledge, Houwang Li and Wenming Zou [16] discussed the normalized solutions for quasilinear problem (1.4) with $f(u) = |u|^{p-2}u$ satisfies $p > 4 + \frac{4}{N}$, and a > 0.

Motivated by the results above, considering that there are few results on normalized solutions for the quasilinear Schrödinger equation with exponential critical growth, in this paper, we focus on the system (1.4) discussed before, where $a \in (0,1), \lambda \in \mathbb{R}$ and f has an exponential critical growth. We recall that in \mathbb{R}^2 , the natural growth restriction on function f is given by the inequality of Trudinger and Moser [24, 34].

In this paper, we assume that f is a continuous function that satisfies the following conditions:

$$(f_1) \lim_{t\to 0} \frac{|f(t)|}{|t|^{\tau}} = 0, \text{ for some } \tau > 2 + \frac{\alpha}{2};$$

 (f_2) f has γ_0 exponential critical growth, i.e., there exists $\gamma_0 > 0$ such that

$$\lim_{|t| \to +\infty} \frac{|f(t)|}{e^{\gamma t^2}} = \begin{cases} 0, & \text{for } \gamma > \gamma_0, \\ +\infty, & \text{for } 0 < \gamma < \gamma_0, \end{cases}$$

(f₃) There exists a constant $\kappa > 3 + \frac{\alpha}{2}$ such that

$$0 < \kappa F(t) \le t f(t)$$
, for all $t \in \mathbb{R} \setminus \{0\}$, where $F(t) = \int_0^t f(s) ds$;

 (f_4) There exist constants $\sigma > 3 + \frac{\alpha}{2}$ and $\mu > 0$ such that

$$F(t) \ge \mu |t|^{\sigma}$$
, for all $t \in \mathbb{R} \setminus \{0\}$,

Our main result is as follows:

Theorem 1.1. Assume that f satisfies $(f_1) - (f_4)$. If $a^2 < \frac{(2+\alpha)\pi}{\gamma_0}$, then there exists $\mu^* = \mu^*(a) > 0$ such that problem (1.4) admits a couple of normalized solution $(u_a, \lambda_a) \in H^{1,2}(\mathbb{R}^2) \times \mathbb{R}$ of weak solution, u_a is a radially symmetric function, and $\lambda_a < 0$ for all $\mu \ge \mu^*$.

Remark 1.1. A typical example satisfying $(f_1) - (f_4)$ is

$$f(t) = \mu |t|^{p-2} t e^{\alpha_0} |t|^2$$
, for $p > 3 + \frac{\alpha}{2}$, and all $t \in \mathbb{R}$.

The organization of this paper is as follows: in Section 2, we state some preliminary lemmas and perturbation settings. In Section 3, we use the mountain-pass arguments to construct a bounded (PS)sequence. In Section 4, we prove the existence of critical points for perturbation functional. In Section ??, we study the convergence of critical points for the perturbation functional as $\eta \to 0^+$. Section 5 is devoted to the proof of Theorem 1.1.

Notations:

- Write $|u|_p := \left(\int_{\mathbb{R}^2} |u|^p dx\right)^{\frac{1}{p}}$ with $1 \le p < \infty$. Denote $H^{1,2}(\mathbb{R}^2) := \left\{u \in L^2(\mathbb{R}^2) \mid Du \in L^2(\mathbb{R}^2)\right\}$ as the Sobolev space with the norm

$$||u|| := \left(\int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} |u|^2 dx \right)^{\frac{1}{2}}.$$

- Use " \rightarrow " and " \rightarrow " respectively to denotes the strong and weak convergence in the related function space.
- Define $B_R^c := \{ x \in \mathbb{R}^2 : |x| > R \}.$
- Use " $C_1, C_2, C_3 \cdots$ " denotes any positive constants (possibly different).
- Denote $o_n(1)$ a real sequence with $o_n(1) \to 0$ as $n \to +\infty$.

Next, we will introduce the following Gagliardo-Sobolev inequalitity [1]: for any p > 2,

$$|u|_p \le C_p |\nabla u|_2^{\gamma_p} |u|_2^{1-\gamma_p}, \quad \forall u \in H^{1,2}(\mathbb{R}^2),$$
 (1.6)

where

$$\gamma_p := 2\left(\frac{1}{2} - \frac{1}{p}\right).$$

and the following Gagliardo-Nirenberg-type inequality [16]:

$$\int_{\mathbb{R}^2} |u|^p \ dx \le C_p \left(\int_{\mathbb{R}^2} |u|^2 \ dx \right) \left(4 \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 \ dx \right)^{\frac{p-2}{4}}, \quad \forall u \in H^{1,2}(\mathbb{R}^2). \tag{1.7}$$

2 Preliminaries

2.1 Preliminary lemmas

We start our study recalling that by $(f_1) - (f_3)$ and f(t) has critical exponential growth at $+\infty$ with critical exponential γ_0 , then fix $q > 1 + \frac{\alpha}{2}$, $\tau > 2 + \frac{\alpha}{2}$, for any $\varepsilon > 0$ and $\gamma > \gamma_0$ close to γ_0 , there exists a constant C > 0 which depends on ε and μ such that

$$|f(t)| \le \varepsilon |t|^{\tau} + C_{\varepsilon,\mu} |t|^{q-1} (e^{\gamma t^2} - 1) \quad \text{for all } t \in \mathbb{R},$$
(2.1)

and, it is easy to see that

$$|F(t)| \le \varepsilon |t|^{\tau+1} + C_{\varepsilon,\mu}|t|^q (e^{\gamma t^2} - 1) \quad \text{for all } t \in \mathbb{R}.$$

Now, we recall the following version of Trudinger-Moser inequality as stated in [7].

Lemma 2.1. (i) If $\gamma > 0$ and $u \in H^{1,2}(\mathbb{R}^2)$, then

$$\int_{\mathbb{R}^2} (e^{\gamma u^2} - 1) dx < +\infty.$$

(ii) Moreover, if $|\nabla u|_2^2 \le 1$, $|u|_2^2 \le M < +\infty$, and $0 < \gamma < 4\pi$, then, there exists a constant C > 0 which depends only on M and γ , such that

$$\int_{\mathbb{D}^2} (e^{\gamma u^2} - 1) dx \le C_{\gamma, M}.$$

And the Hardy-Littlewood-Sobolev inequality, see

Lemma 2.2. Let $t, r > 1, 0 < \alpha < 2$, with $\frac{1}{t} + \frac{2-\alpha}{2} + \frac{1}{r} = 2$, $f \in L^t(\mathbb{R}^2)$ and $g \in L^r(\mathbb{R}^2)$. Then, there exists a sharp constant C which depends on t, α, r , such that

$$\int_{\mathbb{R}^2} (I_{\alpha} * f) g dx \le C_{t,\alpha,r} |f|_t |g|_r.$$

According to Lemma 2.2, we arrive at

$$\int_{\mathbb{R}^2} (I_\alpha * F(u)) F(u) dx$$

is well-defined if $F(u) \in L^t(\mathbb{R}^2)$ for t > 1 given by

$$\frac{2}{t} + \frac{2-\alpha}{2} = 2.$$

This implies that we must require

$$F(u) \in L^{\frac{4}{2+\alpha}}(\mathbb{R}^2).$$

We also need the following inequality, which will be used in the following lemmas.

$$(e^s - 1)^t \le e^{ts} - 1$$
, for $t > 1$ and $s \ge 0$.

Lemma 2.3. Assume that $\{u_n\} \subset \widetilde{S}(a)$ is bounded and satisfies

$$\limsup_{n \to \infty} |\nabla u_n|_2^2 \in \left(0, \frac{(2+\alpha)\pi}{\gamma} - a^2\right).$$

Then for $\gamma > \gamma_0$ close to γ_0 , the sequence $\{e^{\gamma |u_n|^2} - 1\}$ is bounded in $L^t(\mathbb{R}^2)$ provided t > 1 close to 1. Proof. Write

$$\beta := \limsup_{n \to \infty} ||u_n||_2^2.$$

Under the assumptions, we have that $\beta \in \left(0, \frac{(2+\alpha)\pi}{\gamma_0}\right)$. Then we can find some $\eta > 0$ such that $\beta < \frac{(2+\alpha)\pi}{\gamma_0 + \eta}$. Without loss of generality, we may assume that

$$||u_n||_2^2 < \frac{(2+\alpha)\pi}{\gamma_0 + \eta}, \quad \forall n \in \mathbb{N}.$$

Since $\gamma > \gamma_0$ close to γ_0 , we can write it as $\gamma = \gamma_0 + \xi$. Letting $t \in (1, 1+\xi)$ with $\xi \in (0, \min\{\frac{\eta}{\gamma_0 + 2}, 1\})$, we have

$$\limsup_{n \to \infty} t\gamma ||u_n||_2^2 \le \limsup_{n \to \infty} (1+\xi)(\gamma_0 + \xi)||u_n||_2^2 \le (\eta + \gamma_0)\beta < (2+\alpha)\pi < 4\pi.$$

Noting that $\beta > 0$, and by Lemma2.1, we arrive at

$$\limsup_{n \to \infty} \int_{\mathbb{R}^2} (e^{\gamma |u_n|^2} - 1)^t dx \le \limsup_{n \to \infty} \int_{\mathbb{R}^2} (e^{t\gamma |u_n|^2} - 1) dx$$

$$= \limsup_{n \to \infty} \int_{\mathbb{R}^2} (e^{t\gamma ||u_n||^2 \left(\frac{|u_n|}{||u_n||}\right)^2} - 1) dx$$

$$< +\infty.$$

The lemma is finished.

Corollary 2.1. Assume that $u_n \rightharpoonup u_0$ weakly in $\widetilde{S}(a)$ and $\limsup_{n \to \infty} |\nabla (u_n - u_0)|_2^2 < \frac{(2+\alpha)\pi}{\gamma} - a^2$. Then for $\gamma > \gamma_0$ close to γ_0 , we have that $\{e^{\gamma|u_n|^2} - 1\}$ is bounded in $L^t(\mathbb{R}^2)$ provided t > 1 close to 1.

Proof. We just need to prove that when n large enough, for $\gamma > \gamma_0$ and t > 1 close to 1, it still holds

$$\lim_{n \to \infty} \sup t\gamma ||u_n - u_0||_2^2 < 4\pi.$$
 (2.3)

Let $v_n = u_n - u_0$. By choosing $\omega > 0$ small enough, there exists some $C_\omega > 0$ such that

$$\int_{\mathbb{R}^{2}} (e^{\gamma|u_{n}|^{2}} - 1)^{t} dx \leq \int_{\mathbb{R}^{2}} (e^{t\gamma|u_{n}|^{2}} - 1) dx
= \int_{\mathbb{R}^{2}} (e^{t\gamma|v_{n} + u_{0}|^{2}} - 1) dx
\leq \int_{\mathbb{R}^{2}} (e^{(1+\omega)t\gamma|v_{n}|^{2} + C_{\omega}t\gamma|u_{0}|^{2}} - 1) dx.
= \int_{\mathbb{R}^{2}} (e^{(1+\omega)t\gamma|v_{n}|^{2}} - 1) (e^{C_{\omega}t\gamma|u_{0}|^{2}} - 1) dx + \int_{\mathbb{R}^{2}} e^{(1+\omega)t\gamma|v_{n}|^{2}} dx + \int_{\mathbb{R}^{2}} e^{C_{\omega}t\gamma|u_{0}|^{2}} dx - 2
:= I + II + III + IV.$$

By choosing r > 1 close to 1, according to (2.3), Lemma 2.1 and the Hölder inequality, there exists some C > 0 independent of n such that

$$I := \int_{\mathbb{R}^2} (e^{(1+\omega)t\gamma|v_n|^2} - 1)(e^{C_\omega t\gamma|u_0|^2} - 1)dx$$

$$\leq \left(\int_{\mathbb{R}^2} (e^{(1+\omega)t\gamma|v_n|^2} - 1)^r dx\right)^{\frac{1}{r}} \left(\int_{\mathbb{R}^2} (e^{C_\omega t\gamma|u_0|^2} - 1))^{r'} dx\right)^{\frac{1}{r'}}$$

$$\leq C,$$

$$II := \int_{\mathbb{R}^2} e^{(1+\omega)t\gamma|v_n|^2} dx \le C,$$

$$III := \int_{\mathbb{R}^2} e^{C_{\omega} t \gamma |u_0|^2} dx \le C,$$

where $r' = \frac{r}{r-1}$. Hence, $\{e^{\gamma |u_n|^2} - 1\}$ is bounded in $L^t(\mathbb{R}^2)$

Lemma 2.4. Assume that $\{u_n\}\subset \widetilde{S}(a)$ such that $u_n\rightharpoonup u_0$ weakly in $H^{1,2}_{rad}(\mathbb{R}^2)\cap \widetilde{S}(a)$ and

$$\lim_{n \to \infty} \sup |\nabla u|_2^2 < \frac{2+\alpha}{\gamma} - a^2. \tag{2.4}$$

then we arrive at

$$|u_n|^q (e^{\gamma |u_n|^2} - 1) \to |u|^q (e^{\gamma |u_0|^2} - 1)$$
 in $L^t(\mathbb{R}^2)$

Proof. Setting

$$h_n(x) = e^{\gamma |u_n|^2 - 1}.$$

By (2.4) and Lemma 2.3, we have $\{h_n\}$ is a bounded sequence in $L^t(\mathbb{R}^2)$. By $u_n \rightharpoonup u_0$ in $H^{1,2}_{rad}(\mathbb{R}^2)$, we know that $u_n \to u$ a.e. in \mathbb{R}^2 . Thus, we obtain $h_n(x) = e^{\gamma |u_n|^2 - 1} \to e^{\gamma |u|^2 - 1}$ a.e. in \mathbb{R}^2 . Then, we have

$$h_n \rightharpoonup h = e^{\gamma |u|^2 - 1} \quad in \quad L^t(\mathbb{R}^2).$$
 (2.5)

Now we show that

$$|u_n|^q \to |u|^q \quad in \quad L^{t'}(\mathbb{R}^2),$$
 (2.6)

where $t' = \frac{t}{t-1}$. Then by the embedding $H^{1,2}_{rad}(\mathbb{R}^2) \hookrightarrow L^{qt'}(\mathbb{R}^2)$ is compact, we have

$$u_n \to u \quad L^{qt'}(\mathbb{R}^2).$$

Thus, we get (2.6). Together (2.5) with (2.6), we know

$$|u_n|^q (e^{\gamma |u_n|^2} - 1) \to |u|^q (e^{\gamma |u|^2} - 1) in \quad L^1(\mathbb{R}^2).$$

Then, the proof is complete.

Corollary 2.2. Assume that $(f_1) - (f_3)$ hold, let $\{u_n\} \subset \widetilde{S}(a) \cap H^{1,2}_{rad}(\mathbb{R}^2)$ with

$$\limsup_{n \to \infty} |\nabla u_n|_2^2 < \frac{(2+\alpha)\pi}{\gamma} - a^2.$$

If $u_n \rightharpoonup u$ in $H^{1,2}_{rad}(\mathbb{R}^2)$ and $u_n(x) \rightarrow u(x)$ a.e. in \mathbb{R} , then

$$\int_{\mathbb{R}^2} \left(I_\alpha * F(u_n) \right) f(u_n) \phi dx \to \int_{\mathbb{R}^2} \left(I_\alpha * F(u) \right) f(u) \phi dx, \ as \ n \to \infty,$$

for any $\phi \in C_0^{\infty}(\mathbb{R}^2)$.

Proof. First, we claim that $I_{\alpha} * F(u_n)$ belongs to $L^{\infty}(\mathbb{R}^2)$, indeed, by (2.2), we have

$$|F(u_n)| \le \varepsilon |u_n|^{\tau+1} + K_{\varepsilon}|u_n|^q (e^{\gamma |u_n|^2} - 1).$$

Then

$$\begin{aligned} |I_{\alpha} * F(u_{n})| &= \left| \int_{\mathbb{R}^{2}} \frac{A_{\alpha}}{|x - y|^{2 - \alpha}} F(u_{n}) dx \right| \\ &= \left| \int_{|x - y| \le 1} \frac{A_{\alpha}}{|x - y|^{2 - \alpha}} F(u_{n}) dx \right| + K_{\varepsilon} \left| \int_{|x - y| \ge 1} \frac{A_{\alpha}}{|x - y|^{2 - \alpha}} F(u_{n}) dx \right| \\ &= \int_{|x - y| \le 1} \frac{A_{\alpha} \varepsilon |u_{n}|^{\tau + 1}}{|x - y|^{2 - \alpha}} dx + \int_{|x - y| \le 1} \frac{A_{\alpha} K_{\varepsilon} |u_{n}|^{q} (e^{\gamma |u_{n}|^{2}} - 1)}{|x - y|^{2 - \alpha}} dx + K_{\varepsilon} \int_{|x - y| \ge 1} \frac{A_{\alpha} \varepsilon |u_{n}|^{\tau + 1}}{(x - y)^{2 - \alpha}} dx \\ &+ \int_{|x - y| \ge 1} K_{\varepsilon} |u_{n}|^{q} (e^{\gamma |u_{n}|^{2}} - 1) dx \\ &:= I + II + III + IV. \end{aligned}$$

Choose $\sigma \in (\frac{2}{\alpha}, \frac{4}{2+\alpha})$, by the Hölder inequality, we get

$$I := \int_{|x-y| \le 1} \frac{A_{\alpha} \varepsilon |u_n|^{\tau+1}}{|x-y|^{2-\alpha}} dx$$

$$\le A_{\alpha} \varepsilon \left(\int_{|x-y| \le 1} |u_n|^{(\tau+1)\sigma} dx \right)^{\frac{1}{\sigma}} \left(\int_{|x-y| \le 1} \frac{1}{|x-y|^{(2-\alpha)\sigma'}} dx \right)^{\frac{1}{\sigma'}}$$

$$< C_1.$$

$$II := \int_{|x-y| \le 1} \frac{A_{\alpha} K_{\varepsilon} |u_n|^q (e^{\gamma |u_n|^2} - 1)}{|x-y|^{2-\alpha}} dx$$

$$\le C A_{\alpha} K_{\varepsilon} \left(\int_{|x-y| \le 1} |u_n|^{q\sigma t'} dx \right)^{\frac{1}{t'\sigma}} \left(\int_{|x-y| \le 1} (e^{\gamma \sigma t |u_n|^2} - 1) dx \right)^{\frac{1}{t}}$$

$$< C_2.$$

Choose $\delta \to 0^+, st: q_{1,\delta} = \frac{(\tau+1)(2+\delta)}{\delta+\alpha} > 2$, and t > 1 close to 1, by the Hölder inequality, we get

$$III := \int_{|x-y| \ge 1} \frac{A_{\alpha} \varepsilon |u_n|^{\tau+1}}{(x-y)^{2-\alpha}} dx \le A_{\alpha} \varepsilon \left(\int_{|x-y| \le 1} \frac{1}{|x-y|^{2+\delta}} dx \right)^{\frac{2-\alpha}{2+\delta}} \left(\int_{|x-y| \le 1} |u_n|^{q_1,\delta} dx \right)^{\frac{\delta+\alpha}{\tau+\delta}} < C_3,$$

$$IV := \int_{|x-y| \ge 1} K_{\varepsilon} |u_n|^q (e^{\gamma |u_n|^2} - 1) dx \le K_{\varepsilon} \left(\int_{|x-y| \le 1} |u_n|^{qt'} dx \right)^{\frac{1}{t'}} \left(\int_{|x-y| \le 1} (e^{t\gamma |u_n|^2} - 1) dx \right)^{\frac{1}{t}} < C_4,$$

where $\sigma' = \frac{\sigma}{\sigma} - 1$, $t' = \frac{t}{t-1}$. This prove the claim.

Hence, for any $\phi \in C_0^{\infty}(\mathbb{R}^2)$, we have

$$|(I_{\alpha} * F(u_n))f(u_n)\phi| \le C |f(u_n)| |\phi| \le \varepsilon |u_n|^{\tau} |\phi| + C |u_n|^{q-1} |\phi| (e^{\gamma |u_n|^2} - 1).$$

Let $U = supp \phi$. Then, by Lemma 2.4, we obtain

$$\int_{U} |u_{n}|^{\tau} |\phi| dx \to \int_{U} |u|^{\tau} |\phi| dx, \quad as \, n \to \infty,$$

and

$$\int_{U} |u_{n}|^{q-1} |\phi| (e^{\gamma |u_{n}|^{2}} - 1) dx \to \int_{U} |u|^{q-1} |\phi| (e^{\gamma |u|^{2}} - 1) dx, \quad as \, n \to \infty,$$

Now, applying a variant of the Lebesgue dominated convergence theorem, we can deduce that

$$\int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) \phi dx \to \int_{\mathbb{R}^2} (I_\alpha * F(u)) f(u) \phi dx, \quad as n \to \infty.$$

which completes the proof.

Corollary 2.3. Assume that $(f_1) - (f_3)$ hold, let $\{u_n\} \subset \widetilde{S}(a) \cap H^{1,2}_{rad}(\mathbb{R}^2)$ with

$$\limsup_{n \to \infty} |\nabla u_n|_2^2 < \frac{(2+\alpha)\pi}{\gamma} - a^2.$$

If $u_n \rightharpoonup u$ in $H^{1,2}_{rad}(\mathbb{R}^2)$ and $u_n(x) \rightarrow u(x)$ a.e. in \mathbb{R} , then

$$\int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx \to \int_{\mathbb{R}^2} (I_\alpha * F(u)) F(u) dx,$$

and

$$\int_{\mathbb{R}^2} \left(I_\alpha * F(u_n) \right) f(u_n) u_n \ dx \to \int_{\mathbb{R}^2} \left(I_\alpha * F(u) \right) f(u) u dx.$$

Proof. From Corollary 2.2, we know

$$|I_{\alpha} * F(u_n)| \leqslant C.$$

By (2.2), we have

$$|F(u_n)| \le \varepsilon |u_n|^{\tau+1} + K_{\varepsilon} |u_n|^q (e^{\gamma |u_n|^2} - 1).$$

where $\gamma > \gamma_0, \tau > 2 + \frac{2}{\alpha}, q > 1 + \frac{\alpha}{2}$. Hence, we have

$$|(I_{\alpha} * F(u_n))F(u_n)| \le C|F(u_n)| \le \varepsilon |u_n|^{\tau+1} + C|u_n|^q (e^{\gamma |u_n|^2} - 1).$$

By the compact embedding $H^{1,2}_{rad}(\mathbb{R}^2) \hookrightarrow L^p(\mathbb{R}^2)$ for p > 2, we have

$$u_n \to u$$
 in $L^p(\mathbb{R}^2)$.

Now, applying a variant of the Lesbesgue dominated convergence theorem, we can deduce that

$$\int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx \to \int_{\mathbb{R}^2} (I_\alpha * F(u)) F(u) dx, \quad as \, n \to \infty.$$

A similar argument works to show that

$$\int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n \ dx \to \int_{\mathbb{R}^2} (I_\alpha * F(u)) f(u) u dx, \quad as \, n \to \infty.$$

2.2 Perturbation setting

In order to recover the differentiability, we define for $\eta \in (0,1]$,

$$I_{\eta}(u) := \frac{\eta}{\theta} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + I(u)$$

$$= \frac{1}{2} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx + \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx - \frac{1}{2} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u)) F(u) dx$$
(2.7)

on the space $O := H^{1,\theta}(\mathbb{R}^2) \cap H^{1,2}(\mathbb{R}^2)$, for some fixed θ , satisfying $2 < \theta < 3$. Then O is a reflexive Banach space, and by the Hardy-Littlewood-Sobolev inequality and [16, Lemma A.1], we can know that $I_{\eta} \in C^1(O)$. We will consider I_{η} on the constraint

$$S(a) := \left\{ u \in O : \int_{\mathbb{R}^2} |u|^2 dx = a^2 \right\}. \tag{2.8}$$

Recalling the L^2 -norm preserved transform [15] $\mathcal{H}: H^{1,2}(\mathbb{R}^2) \times \mathbb{R} \to \mathbb{R}$ and

$$\mathcal{H}(u,s)(x) = e^s u(e^s x).$$

Then, we have

$$I_{\eta}(\mathcal{H}(u,s)) := \frac{\eta}{\theta} e^{2(\theta-1)s} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + e^{4s} \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} e^{2s} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx - \frac{1}{2e^{(2+\alpha)s}} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(e^{s}u)) F(e^{s}u) dx.$$
(2.9)

We define

$$P_{\eta}(u) := \frac{d}{ds}|_{s=0}I_{\eta}(\mathcal{H}(u,s))$$

$$= \eta \frac{2(\theta-1)}{\theta} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + 4 \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx + \frac{2+\alpha}{2} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u)) F(u) dx$$

$$- \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u)) f(u) u dx, \qquad (2.11)$$

and $P_{\eta} \in C^{1}(O)$, then we define a manifold

$$\mathcal{P}_{\eta}(a) := \{ u \in \mathcal{S}(a) : P_{\eta}(u) = 0 \}. \tag{2.12}$$

We have the following results.

Lemma 2.5. Any critical point u of $I_{\eta}|_{\mathcal{S}(a)}$ is contained in $\mathcal{P}_{\eta}(a)$.

Proof. By [6, Lemma 3], there exists a $\lambda \in \mathbb{R}$ such that

$$I'_{\eta}(u) + \lambda u = 0 \quad \text{in} \quad O^*.$$
 (2.13)

On the one hand, using (2.13), we obtain

$$\eta \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx + 4 \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx + \int_{\mathbb{R}^2} |\nabla u|^2 dx + \lambda \int_{\mathbb{R}^2} |u|^2 dx - 2 \int_{\mathbb{R}^2} (I_{\alpha} * F(u)) f(u) u dx = 0. \quad (2.14)$$

On the other hand, testing (2.13), for more details see [5, Proposition 1], we obtain

$$\eta \frac{\theta - 2}{2\theta} \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx - \frac{1}{2} \lambda \int_{\mathbb{R}^2} |u|^2 dx + \frac{2 + \alpha}{4} \int_{\mathbb{R}^2} (I_{\alpha} * F(u)) F(u) dx + \frac{1}{2} \int_{\mathbb{R}^2} (I_{\alpha} * F(u)) f(u) u = 0. \quad (2.15)$$

Combining (2.14) and (2.15), we can get $P_{\eta}(u) = 0$, so $u \in \mathcal{P}_{\eta}(a)$.

Lemma 2.6. The following statements hold: if $\sup_{n\geq 1} I_{\eta}(u_n) < +\infty$ for $u_n \in \mathcal{P}_{\eta}(a)$, then

$$\sup_{n\geq 1} \max\left\{\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx, \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx, \int_{\mathbb{R}^2} |\nabla u_n|^2 dx\right\} < +\infty.$$

Proof. For any $u \in \mathcal{P}_{\eta}(a)$, exists $z \in \mathbb{R}$ with $\frac{1}{2k-2-\alpha} < z < \frac{1}{4}$, $k > 3 + \frac{2}{\alpha}$, by (f_3) we have

$$\begin{split} I_{\eta}(u) &= I_{\eta}(u) - zP_{\eta}(u) \\ &= \eta(\frac{1 - 2(\theta - 1)z}{\theta}) \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + (1 - 4z) \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + (\frac{1 - 2z}{2}) \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx \\ &- \frac{(2 + \alpha)z + 1}{2} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u))F(u) dx + z \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u))F(u) dx \\ &\geq \eta(\frac{1 - 2(\theta - 1)z}{\theta}) \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + (1 - 4z) \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx - (\frac{(2 + \alpha)z + 1}{2}) \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u))F(u) dx \\ &+ kz \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u))F(u) dx \\ &= \eta(\frac{1 - 2(\theta - 1)z}{\theta}) \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + (1 - 4z) \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + (\frac{1 - 2z}{2}) \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx \\ &+ \left(kz - (\frac{(2 + \alpha)z + 1}{2})\right) \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u))F(u) dx \end{split}$$

As for $\frac{1-(2\theta-2)z}{\theta} > 0$, 1-4z > 0, $\frac{1-2z}{2} > 0$, $kz - (\frac{(2+\alpha)z+1}{2}) > 0$, so the conclusion has finished.

3 The minimax approach

In this section, we will prove that $I_{\eta}(\mathcal{H}(u,s))$ on $S(a) \times \mathbb{R}$ possesses a kind of mountain-pass geometrical structure.

Lemma 3.1. For any $0 < \eta \le 1$ and $u \in S(a)$ be arbitrary but fixed, the following statements hold:

- (i) $|\nabla \mathcal{H}(u,s)|_2 \to 0^+$ and $I_{\eta}(\mathcal{H}(u,s)) \to 0^+$ as $s \to -\infty$;
- (ii) $|\nabla \mathcal{H}(u,s)|_2 \to +\infty$ and $I_n(\mathcal{H}(u,s)) \to -\infty$ as $s \to +\infty$.

Proof. By a straightforward calculation, it follows that

$$\int_{\mathbb{R}^2} |\mathcal{H}(u,s)(x)|^2 dx = a^2, \quad \int_{\mathbb{R}^2} |\mathcal{H}(u,s)(x)|^{\varsigma} dx = e^{(\varsigma-2)s} \int_{\mathbb{R}^2} |u(x)|^{\varsigma} dx \quad \text{for all } \varsigma > 2,$$

$$\int_{\mathbb{R}^2} |\nabla \mathcal{H}(u,s)(x)|^2 dx = e^{2s} \int_{\mathbb{R}^2} |\nabla u(x)|^2 dx, \quad \int_{\mathbb{R}^2} |\nabla \mathcal{H}(u,s)(x)|^{\theta} dx = e^{(2\theta-2)s} \int_{\mathbb{R}^2} |\nabla u(x)|^{\theta} dx,$$

$$\int_{\mathbb{R}^2} |\mathcal{H}(u,s)(x)|^2 |\nabla \mathcal{H}(u,s)(x)|^2 dx = e^{4s} \int_{\mathbb{R}^2} |u(x)|^2 |\nabla u(x)|^2 dx.$$

From the above equalities, fixing $\varsigma > 2$, as $s \to -\infty$, it follows that

$$|\mathcal{H}(u,s)|_{\varsigma}^{\varsigma} \to 0^{+}, \quad |\nabla \mathcal{H}(u,s)|_{2}^{2} \to 0^{+}, \quad |\nabla \mathcal{H}(u,s)|_{\theta}^{\theta} \to 0^{+}, \quad \int_{\mathbb{R}^{2}} |\mathcal{H}(u,s)|^{2} |\nabla \mathcal{H}(u,s)|^{2} dx \to 0^{+}.$$

By (2.2), we have

$$|F(\mathcal{H}(u,s))| \le \varepsilon |\mathcal{H}(u,s)|^{\tau+1} + K_{\varepsilon,\mu} |\mathcal{H}(u,s)|^q (e^{\gamma |\mathcal{H}(u,s)|^2} - 1), \tag{3.1}$$

For all $\gamma ||\mathcal{H}(u,s)||^2 < (2+\alpha)\pi$, by Lemma2.1, we have

$$\int_{\mathbb{R}^2} \left(e^{\gamma |\mathcal{H}(u,s)|^2} - 1 \right) dx = \int_{\mathbb{R}^2} \left(e^{\gamma |\mathcal{H}(u,s)|^2 \left(\frac{|\mathcal{H}(u,s)|}{|\mathcal{H}(u,s)|} \right)^2} - 1 \right) dx \le C.$$

Hence, using Hardy-Littlewood-Sobolev inequality, Minkowski inequality and the Hölder's inequality, we deduce that

$$\left(\int_{\mathbb{R}^{2}} |F(\mathcal{H}(u,s))|^{\frac{4}{2+\alpha}} dx\right)^{\frac{2+\alpha}{4}} \leq \left(\int_{\mathbb{R}^{2}} [\varepsilon|\mathcal{H}(u,s)|^{\tau+1} dx + K_{\varepsilon,\mu}|\mathcal{H}(u,s)|^{q} (e^{\gamma|\mathcal{H}(u,s)|^{2}} - 1)]^{\frac{4}{2+\alpha}} dx\right)^{\frac{2+\alpha}{4}} \\
\leq \left(\int_{\mathbb{R}^{2}} \varepsilon|\mathcal{H}(u,s)|^{\frac{4(\tau+1)}{2+\alpha}} dx\right)^{\frac{2+\alpha}{4}} \\
+ K_{\varepsilon,\mu} \left(\int_{\mathbb{R}^{2}} |\mathcal{H}(u,s)|^{\frac{4qt'}{2+\alpha}} dx\right)^{\frac{2+\alpha}{4t'}} \left(\int_{\mathbb{R}^{2}} e^{\frac{4t\gamma}{2+\alpha}|\mathcal{H}(u,s)|^{2}} - 1 dx\right)^{\frac{2+\alpha}{4t}},$$

where t, t' > 1 satisfying $\frac{1}{t} + \frac{1}{t'} = 1$. Then, there exists t > 1 close to 1 such that

$$t\gamma ||\mathcal{H}(u,s)||^2 < (2+\alpha)\pi,$$

which implies that

$$\left(\int_{\mathbb{R}^2} e^{\frac{4t\gamma}{2+\alpha}|\mathcal{H}(u,s)|^2} - 1dx\right)^{\frac{2+\alpha}{4t}} \le C. \tag{3.2}$$

Note that, by (3.2) and the Hölder's inequality, we arrive at

$$\int_{\mathbb{R}^{2}} \left(I_{\alpha} * F(\mathcal{H}(u,s)) \right) F(\mathcal{H}(u,s) dx \leq |F(\mathcal{H}(u,s))|_{\frac{4}{2+\alpha}} |F(\mathcal{H}(u,s))|_{\frac{4}{2+\alpha}}$$

$$\leq \left(\varepsilon \left| \mathcal{H}(u,s) \right|_{\frac{4(\tau+1)}{2+\alpha}}^{\frac{\tau+1}{2+\alpha}} + C \left| \mathcal{H}(u,s) \right|_{\frac{4qt'}{2+\alpha}}^{q} \right)^{2}.$$

Thus, we conclude that

$$\int_{\mathbb{R}^2} \left(I_\alpha * F(\mathcal{H}(u,s)) \right) F(\mathcal{H}(u,s) dx \le \left(\varepsilon e^{\frac{2\tau - \alpha}{2}s} |u|_{\frac{4(\tau + 1)}{2 + \alpha}}^{\frac{\tau + 1}{4(\tau + 1)}} + C e^{\frac{4qt' - 4 - 2\alpha}{4t'}s} |u|_{\frac{4qt'}{2 + \alpha}}^q \right)^2.$$

And then, we have

$$\begin{split} I_{\eta}(\mathcal{H}(u,s)) & \geq \frac{\eta}{\theta} e^{2(\theta-1)s} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + e^{4s} \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} e^{2s} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx \\ & - \frac{1}{2e^{(2+\alpha)s}} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(e^{s}u)) F(e^{s}u) dx \\ & \geq \frac{\eta}{\theta} e^{2(\theta-1)s} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + e^{4s} \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} e^{2s} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx \\ & - \left(\varepsilon e^{\frac{2\tau-\alpha}{2}s} |u|_{\frac{4(\tau+1)}{2+\alpha}}^{\tau+1} + C e^{\frac{(4qt'-4-2\alpha)s}{4t'}} |u|_{\frac{4qt'}{2+\alpha}}^{q} \right)^{2}. \\ & := I_{\eta,1}(\mathcal{H}(u,s)) \end{split}$$

Thus, by $\tau > 2 + \frac{\alpha}{2}, q > 1 + \frac{\alpha}{2}$, we know that

$$I_{\eta,1}(\mathcal{H}(u,s)) \to 0^+$$
 as $s \to -\infty$,

and

$$I_{\eta,1}(\mathcal{H}(u,s)) \to -\infty$$
 as $s \to +\infty$.

On the other hand, we define

$$g(z) = \int_{\mathbb{R}^2} (I_\alpha * F(z)) F(z) dx,$$

by (2.9),

$$I_{\eta}(\mathcal{H}(u,s)) := \frac{\eta}{\theta} e^{2(\theta-1)s} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + e^{4s} \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} e^{2s} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx - \frac{1}{2e^{(2+\alpha)s}} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(e^{s}u)) F(e^{s}u) dx.$$

Set

$$w(t) = g(\frac{tu}{||u||}),$$

where $t = e^s$. By f(3), we know

$$\frac{w'(t)}{w(t)} \ge \frac{2k}{t}$$

then, by integral operation, we obtain

$$g(tu) \ge g(\frac{u}{||u||})s^{2k}||u||^{2k}.$$

Therefore, we have

$$I_{\eta}(\mathcal{H}(u,s)) \le C_5 e^{2s} + C_6 e^{2(\theta-1)s} + C_7 e^{4s} - C_8 e^{(2k-(2+\alpha))s}$$

:= $I_{\eta,2}(\mathcal{H}(u,s))$

Thus, by $2 < 2(\theta - 1) < 4 < 2k - (2 + \alpha)$, we know that

$$I_{n,2}(\mathcal{H}(u,s)) \to 0^+$$
 as $s \to -\infty$,

and

$$I_{\eta,2}(\mathcal{H}(u,s)) \to -\infty$$
 as $s \to +\infty$.

Moreover, the inequality below also yields that

$$I_{\eta}(\mathcal{H}(u,s)) \to 0^+$$
 as $s \to -\infty$,

$$I_{\eta}(\mathcal{H}(u,s)) \to -\infty$$
 as $s \to +\infty$.

To recover the compactness, we shall study I_{η} on the radial space:

$$S_r(a) := S(a) \cap O_r, \quad O_r := H_{rad}^{1,\theta}(\mathbb{R}^2) \cap H_{rad}^{1,2}(\mathbb{R}^2).$$

Lemma 3.2. There exists $K(a, \mu) > 0$ small enough such that

$$0 < \sup_{u \in \mathcal{A}} I_{\eta}(u) < \inf_{u \in \mathcal{B}} I_{\eta}(u)$$
(3.3)

with

$$\mathcal{A} = \left\{ u \in S_r(a), \eta \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx + \int_{\mathbb{R}^2} |\nabla u|^2 dx \le K(a, \mu) \right\},$$

and

$$\mathcal{B} = \left\{ u \in S_r(a), \eta \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx + \int_{\mathbb{R}^2} |\nabla u|^2 dx = 3K(a, \mu) \right\}.$$

Moreover, $K(a, \mu) \to 0$ when $\mu \to \infty$.

Proof. Firstly, for $u \in \mathcal{A}$, and $v \in \mathcal{B}$, we have the following estimations

$$\frac{\eta}{\theta} \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla u|^2 dx$$

$$\leq \max \left\{ \frac{1}{\theta}, \frac{1}{2}, 1 \right\} \left(\eta \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx + \int_{\mathbb{R}^2} |\nabla u|^2 dx \right)$$

$$\leq K(a, \mu), \tag{3.4}$$

and

$$\frac{\eta}{\theta} \int_{\mathbb{R}^2} |\nabla v|^{\theta} dx + \int_{\mathbb{R}^2} |v|^2 |\nabla v|^2 dx + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla v|^2 dx
\geq \min \left\{ \frac{1}{\theta}, \frac{1}{2}, 1 \right\} \left(\eta \int_{\mathbb{R}^2} |\nabla v|^{\theta} dx + \int_{\mathbb{R}^2} |v|^2 |\nabla v|^2 dx + \int_{\mathbb{R}^2} |\nabla v|^2 dx \right)
= \frac{3}{\theta} K(a, \mu).$$
(3.5)

Now, let $K(a, \mu) < \frac{(2+\alpha)\pi}{3\gamma} - \frac{a^2}{3}$. Thus, we have

$$||v||^2 = |\nabla v|_2^2 + |v|_2^2 < 3K(a,\mu) + a^2 < \frac{(2+\alpha)\pi}{\gamma}.$$

Then, similar as Lemma 3.1, we obtain

$$\int_{\mathbb{R}^2} (I_\alpha * F(v)) F(v) dx \le \left(\varepsilon |v|_{\frac{4(\tau+1)}{2+\alpha}}^{\frac{\tau+1}{4(\tau+1)}} + C|v|_{\frac{4qt'}{2+\alpha}}^q \right)^2$$

$$\le \varepsilon |v|_{\frac{4(\tau+1)}{2+\alpha}}^{\frac{2(\tau+1)}{2+\alpha}} + C|v|_{\frac{4qt'}{2+\alpha}}^{\frac{2q}{2+\alpha}}$$

where $\tau > 2 + \frac{\alpha}{2}, q > 1 + \frac{\alpha}{2}$. By the Gagliardo-Sobolev inequality (1.6), we have

$$\int_{\mathbb{R}^2} (I_\alpha * F(v)) F(v) dx \le C |\nabla v|_2^{2\tau - \alpha} a^{\frac{2+\alpha}{2}} + C |\nabla v|_2^{\frac{2qt' - 2 - \alpha}{t'}} a^{\frac{2+\alpha}{2t'}}.$$

From (f_3) , we have $(I_{\alpha} * F(u))F(u) > 0$ for any $u \in H^{1,2}(\mathbb{R})$, then, we have

$$I_{\eta}(v) - I_{\eta}(u) \ge \left(\frac{\eta}{\theta} \int_{\mathbb{R}^2} |\nabla v|^{\theta} dx + \int_{\mathbb{R}^2} |v|^2 |\nabla v|^2 dx + \frac{1}{2} \int_{\mathbb{R}^2} |\nabla v|^2 dx\right)$$

$$-\left(\frac{\eta}{\theta}\int_{\mathbb{R}^{2}}|\nabla u|^{\theta}dx + \int_{\mathbb{R}^{2}}|u|^{2}|\nabla u|^{2}dx + \frac{1}{2}\int_{\mathbb{R}^{2}}|\nabla u|^{2}dx\right) - \frac{1}{2}\int_{\mathbb{R}^{2}}(I_{\alpha}*F(v))F(v)dx$$

$$\geq \frac{3}{\theta}K(a,\mu) - K(a,\mu) - C\left(K(a,\mu)\right)^{\frac{2\tau-\alpha}{2}}a^{\frac{2+\alpha}{2}} - C\left(K(a,\mu)\right)^{\frac{2qt'-2-\alpha}{2t'}}a^{\frac{2+\alpha}{2t'}}$$

$$= \frac{3-\theta}{\theta}K(a,\mu) - C\left(K(a,\mu)\right)^{1+\frac{2\tau-(2+\alpha)}{2}} - C\left(K(a,\mu)\right)^{1+\frac{2qt'-2-\alpha-2}{2t'}}$$

Since $\tau > 2 + \frac{\alpha}{2}$, $2 < \theta < 3$ and t' > 0 with $\frac{2qt' - 2 - \alpha}{2t'} > 1$ and $\frac{3 - \theta}{\theta} > 0$, fixing

$$K(a,\mu) = min \left\{ \frac{(2+\alpha)\pi}{3\gamma_0} - \frac{a^2}{3}, \left[\frac{3-\theta}{C\theta} \right]^{\frac{2}{2\tau - (2+\alpha)}}, \left[\frac{3-\theta}{C\theta} \right]^{\frac{2t'}{2qt' - 2 - 2t' - \alpha}} \right\}, \tag{3.6}$$

and so,

$$I_{\eta}(v) - I_{\eta}(u) \ge \frac{3 - \theta}{2\theta} K(a, \mu) > 0,$$

which shows the desired result. Finally, in order to prove the limit $K(a,\mu) \to 0$ when $\mu \to \infty$, fix $u_0 \in S_r(a)$ with $\eta \int_{\mathbb{R}^2} |\nabla u_0|^{\theta} dx + \int_{\mathbb{R}^2} |u_0|^2 |\nabla u_0|^2 dx + \int_{\mathbb{R}^2} |\nabla u_0|^2 dx \le K(a,\mu)$. Then, (f_4) together with Lemma3.1 ensure that

$$\frac{\mu^2}{2} \int_{\mathbb{R}^2} (I_\alpha * |u_0|^\sigma) |u_0|^\sigma dx \le \int_{\mathbb{R}^2} (I_\alpha * F(u_0)) F(u_0) dx \le C |u_0|_{\frac{4(\tau+1)}{2+\alpha}}^{2(\tau+1)} + C |u_0|_{\frac{4qt'}{2+\alpha}}^{2q}.$$

Therefore, we must have $C \to \infty$ when $\mu \to \infty$, and so $C \to \infty$ when $\mu \to \infty$. This limit together with (3.6) show that $K(a,\mu) \to 0$ when $\mu \to \infty$.

Similar discussion as the last lemma, we have the following Corollary.

Corollary 3.1. For $K(a,\mu) > 0$ given in (3.6), there holds that $I_{\eta}(u) > 0$, for all $u \in S_r(a)$ with $\eta \int_{\mathbb{T}^2} |\nabla u|^{\theta} dx + \int_{\mathbb{T}^2} |u|^2 |\nabla u|^2 dx + \int_{\mathbb{T}^2} |\nabla u|^2 dx \leq K(a,\mu)$. Moreover,

$$I_{\eta}^{**} = \inf \left\{ I_{\eta}(u) : u \in S_{r}(a) \text{ and } \eta \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx = \frac{K(a, \mu)}{3} \right\} > 0.$$

Proof. By Lemma 3.1 and the Gagliardo-Nirenberg-type inequality, we have

$$\begin{split} I_{\eta}(u) &= \frac{\eta}{\theta} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx - \frac{1}{2} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u)) F(u) dx \\ &\geq \frac{\eta}{\theta} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx - C \left(\int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx \right)^{\frac{2\tau - \alpha}{4}} \\ &- C \left(\int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx \right)^{\frac{2qt' - 2 - \alpha}{4t'}}, \end{split}$$

where $\frac{2\tau-\alpha}{4} > 1$ and $\frac{2qt'-2-\alpha}{2t'} > 1$. For any $u \in \partial \mathcal{A}(K(a,\mu),a)$,

$$\partial \mathcal{A}(K(a,\mu),a) := \left\{ u \in S_r(a) : \eta \int_{\mathbb{R}^2} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^2} |u|^2 |\nabla u|^2 dx + \int_{\mathbb{R}^2} |\nabla u|^2 dx = K(a,\mu) \right\}.$$

For a smaller $0 < \rho < K(a, \mu)$, we can get

$$\inf_{\partial \mathcal{A}(\rho,a)} I_{\eta}(u) \ge \frac{\eta}{\theta} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx + \frac{1}{2} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx - C \left(\int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx \right)^{\frac{2\tau - \alpha}{4}} \\
- C \left(\int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx \right)^{\frac{2qt' - 2 - \alpha}{4t'}}, \\
\ge \frac{\eta}{\theta} \int_{\mathbb{R}^{2}} |\nabla u|^{\theta} dx + \frac{1}{2} \int_{\mathbb{R}^{2}} |\nabla u|^{2} dx + C \int_{\mathbb{R}^{2}} |u|^{2} |\nabla u|^{2} dx \\
\ge C\rho > 0,$$

for $K(a, \mu) > 0$ small enough. the proof is completed.

In what follows, we fix $u_0 \in S_r(a)$ and apply Lemma 3.1, Lemma 3.2 and Corollary 3.1 to get two numbers $s_1 = s_1(u_0, a, \mu) < 0$, and $s_2 = s_2(u_0, a, \mu) > 0$, the functions $u_{1,\mu} = \mathcal{H}(u_0, s_1)$ and $u_{2,\mu} = \mathcal{H}(u_0, s_2)$ satisfy

$$\eta \int_{\mathbb{R}^2} |\nabla u_{1,\mu}|^{\theta} dx + \int_{\mathbb{R}^2} |u_{1,\mu}|^2 |\nabla u_{1,\mu}|^2 dx + \int_{\mathbb{R}^2} |\nabla u_{1,\mu}|^2 dx < \frac{K(a,\mu)}{3} \quad with \quad I_{\eta}(u_{1,\mu}) > 0,$$

and

$$\eta \int_{\mathbb{R}^2} |\nabla u_{2,\mu}|^{\theta} dx + \int_{\mathbb{R}^2} |u_{2,\mu}|^2 |\nabla u_{2,\mu}|^2 dx + \int_{\mathbb{R}^2} |\nabla u_{2,\mu}|^2 dx > 3K(a,\mu) \quad with \quad I_{\eta}(u_{2,\mu}) < 0.$$

Now, following the idea from Jeanjean [15], we fix the following mountain pass level given by

$$\gamma_{\mu}(a) := \inf_{h \in \Gamma} \max_{t \in [0,1]} I_{\eta}(h(t)),$$

where

$$\Gamma = \{ h \in C([0,1], S_r(a)) : \eta \int_{\mathbb{R}^2} |\nabla h(0)|^{\theta} dx + \int_{\mathbb{R}^2} |h(0)|^2 |\nabla h(0)|^2 dx + \int_{\mathbb{R}^2} |\nabla h(0)|^2 dx < \frac{K(a,\mu)}{3},$$

$$I_{\eta}(h(1)) < 0 \}.$$

From Corollary 3.1, there exists $t_0 \in (0,1)$ such that

$$\max_{t \in [0,1]} I_{\eta}(h(t)) \ge I_{\eta}(h(t_0)) \ge I_{\eta}^{**} > 0,$$

where I_{η}^{**} was given in Corollary 3.1. Then we obtain that

$$\gamma_{\mu}(a) \ge I_{\eta}^{**} > 0.$$

Lemma 3.3. There holds $\lim_{\mu\to+\infty} \gamma_{\mu}(a) = 0$.

Proof. In what follows, we set the path $h_0(t) = \mathcal{H}(u_0, (1-t)s_1 + ts_2) \in \Gamma$. Then, by (f_4) ,

$$\gamma_{\mu}(a) \leq \max_{t \in [0,1]} I_{\eta}(h_{0}(t))
\leq \max_{t \in [0,1]} \left\{ \frac{\eta}{\theta} \int_{\mathbb{R}^{2}} |\nabla h_{0}(t)|^{\theta} dx + \frac{1}{2} \int_{\mathbb{R}^{2}} |\nabla h_{0}(t)|^{2} dx + \int_{\mathbb{R}^{2}} |h_{0}(t)|^{2} |\nabla h_{0}(t)|^{2} dx \right\}$$

$$\begin{split} &-\frac{\mu^2}{2} \int_{\mathbb{R}^2} \left(I_\alpha * |h_0(t)|^\sigma \right) |h_0(t)|^\sigma dx \bigg\} \\ &= \max_{t \in [0,1]} \left\{ \frac{\eta}{\theta} r^{2\theta - 2} \int_{\mathbb{R}^2} |\nabla u_0|^\theta dx + \frac{1}{2} r^2 \int_{\mathbb{R}^2} |\nabla u_0|^2 dx + r^4 \int_{\mathbb{R}^2} |u_0|^2 |\nabla u_0|^2 dx \right. \\ &\left. - \frac{\mu^2}{2} r^{2\sigma - (2+\alpha)} \int_{\mathbb{R}^2} \left(I_\alpha * |u_0|^\sigma \right) |u_0|^\sigma dx \bigg\} \\ &:= \max_{t \in [0,1]} g_0(r), \end{split}$$

where $r := e^{(1-t)s_1 + ts_2}$. Case 1. If 0 < r < 1, we have

$$g_{0}(r) \leq r^{2} \max \left\{ \frac{1}{\theta}, \frac{1}{2}, 1 \right\} \left(\eta \int_{\mathbb{R}^{2}} |\nabla u_{0}|^{\theta} dx + \int_{\mathbb{R}^{2}} |\nabla u_{0}|^{2} dx + \int_{\mathbb{R}^{2}} |u_{0}|^{2} |\nabla u_{0}|^{2} dx \right)$$

$$- \frac{\mu^{2}}{2} r^{2\sigma - (2+\alpha)} \int_{\mathbb{R}^{2}} (I_{\alpha} * |u_{0}|^{\sigma}) |u_{0}|^{\sigma} dx$$

$$= r^{2} \left(\eta \int_{\mathbb{R}^{2}} |\nabla u_{0}|^{\theta} dx + \int_{\mathbb{R}^{2}} |\nabla u_{0}|^{2} dx + \int_{\mathbb{R}^{2}} |u_{0}|^{2} |\nabla u_{0}|^{2} dx \right)$$

$$- \frac{\mu^{2}}{2} r^{2\sigma - (2+\alpha)} \int_{\mathbb{R}^{2}} (I_{\alpha} * |u_{0}|^{\sigma}) |u_{0}|^{\sigma} dx$$

$$:= g_{1}(r).$$

It is not difficult to check that g_1 has a unique critical point \tilde{r} on $(0, +\infty)$, which is a global maximum point at positive level.

$$\widetilde{r} = \left[\frac{4 \left(\eta \int_{\mathbb{R}^2} |\nabla u_0|^{\theta} dx + \int_{\mathbb{R}^2} |\nabla u_0|^2 dx + \int_{\mathbb{R}^2} |u_0|^2 |\nabla u_0|^2 dx \right)}{\mu^2 (2\sigma - (2+\alpha)) \int_{\mathbb{R}^2} \left(I_{\alpha} * |u_0|^{\sigma} \right) |u_0|^{\sigma} dx} \right]^{\frac{1}{2\sigma - 4 - \alpha}} > 0,$$

and so.

$$\gamma_{\mu}(a) \le C \left(\frac{1}{\mu}\right)^{\frac{2}{2\sigma - 4 - \alpha}} \to 0, \text{ as } \mu \to +\infty.$$

Case 2. If $r \geq 1$, we have

$$g_0(r) \le r^4 \left(\eta \int_{\mathbb{R}^2} |\nabla u_0|^{\theta} dx + \int_{\mathbb{R}^2} |\nabla u_0|^2 dx + \int_{\mathbb{R}^2} |u_0|^2 |\nabla u_0|^2 dx \right) - \frac{\mu^2}{2} r^{2\sigma - (2+\alpha)} \int_{\mathbb{R}^2} \left(I_{\alpha} * |u_0|^{\sigma} \right) |u_0|^{\sigma} dx.$$

Discussed as before, we have

$$\gamma_{\mu}(a) \le C \left(\frac{1}{\mu}\right)^{\frac{2}{2\sigma - 6 - \alpha}} \to 0, \text{ as } \mu \to +\infty.$$

Hence,

$$\gamma_{\mu}(a) \leq \min \left\{ C\left(\frac{1}{\mu}\right)^{\frac{2}{2\sigma - 4 - \alpha}}, C\left(\frac{1}{\mu}\right)^{\frac{2}{2\sigma - 6 - \alpha}} \right\} \to 0, \text{ as } \mu \to +\infty,$$

for some C > 0 (possibly different) that do not depend on $\mu > 0$.

To find a Palais-Smale sequence, we consider an auxiliary functional

$$\widetilde{I}_{\eta}(s,u) := I_{\eta}(\mathcal{H}(u,s)) = I_{\eta}(h(t)) : \mathbb{R} \times S_{r}(a) \to \mathbb{R}.$$
(3.7)

Notice that \widetilde{I}_{η} is of class C^1 , by the symmetric critical point principle [25], a Palais-Smale sequence for $\widetilde{I}_{\eta}|_{\mathbb{R}\times S_r(a)}$ is also a Palais-Smale sequence for $\widetilde{I}_{\eta}|_{\mathbb{R}\times S(a)}$. Denoting the closed sublevel set by

$$I_{\eta}^{c} = \left\{ u \in S(a) : I_{\eta}(u) \leq c \right\},\,$$

we also define

$$\sigma_{\eta}(a) := \inf_{\widetilde{h} \in \Gamma_{\eta}} \max_{t \in [0,1]} \widetilde{I}_{\eta}(\widetilde{h}(t)),$$

where

$$\Gamma_{\eta} := \left\{ \widetilde{h} = (\gamma, \beta) \in C([0, 1], \mathbb{R} \times S_r(a)) : \eta \int_{\mathbb{R}^2} |\nabla \beta(0)|^{\theta} dx + \int_{\mathbb{R}^2} |\beta(0)|^2 |\nabla \beta(0)|^2 dx + \int_{\mathbb{R}^2} |\nabla \beta(0)|^2 dx \right.$$

$$< \frac{K(a, \mu)}{3}, I_{\eta}(\beta(1)) < 0, \gamma(0) = 0, \gamma(1) = 0 \right\},$$

Obviously, it holds that $\sigma_{\eta}(a) = \gamma_{\mu}(a)$.

The same discussed as in [16, Lemma 3.6], taking a minimizing sequence $\{\widetilde{h}_n = (0, \beta_n)\} \subset \Gamma_{\eta}$ with $\beta_n \geq 0$ a.e. in \mathbb{R}^2 , there exists a Palais-Smale sequence $\{(s_n, w_n)\} \subset \mathbb{R} \times S_r(a)$ for $\widetilde{I}_{\eta}|_{\mathbb{R} \times S_r(a)}$ at level $\sigma_{\eta}(a)$. Let $u_n = \mathcal{H}(w_n, s_n)$, we have

$$-\Delta u_n - u_n \Delta u_n^2 + \lambda_n u_n = (I_\alpha * F(u_n)) F(u_n) + o_n(1), \quad in \ O^*,$$

$$I_n(u_n) \to \sigma_n(a) = \gamma_\mu(a), \quad \text{as} \quad n \to +\infty,$$

$$(3.8)$$

with the additional property that

$$|s_n| + dist_O(w_n, \beta_n([0, 1])) \to 0$$
, as $n \to \infty$,

Moreover, for some sequence $\{\lambda_n\} \subset \mathbb{R}$, and

$$P_{\eta}(u_n) = \eta \frac{2(\theta - 1)}{\theta} \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx$$

$$+ \frac{2 + \alpha}{2} \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) F(u_n) dx - \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) f(u_n) u_n dx$$

$$\to 0, \quad \text{as} \quad n \to +\infty.$$

$$(3.9)$$

From Lemma 2.6, we know that $\{u_n\}$ is bounded in O_r , and so, the number λ_n must satisfy the equality below

$$\lambda_n = \frac{1}{a^2} \left\{ -\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) f(u_n) u_n dx \right\} + o_n(1).$$

Lemma 3.4. There holds

$$\lim_{n \to +\infty} \sup_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx \le \frac{4}{k - 3 - \frac{\alpha}{2}} \gamma_\mu(a).$$

Proof. Using the fact that $I_{\eta}(u_n) = \gamma_{\mu}(a) + o_n(1)$ and $P_{\eta}(u_n) = o_n(1)$, it follows that

$$\eta \frac{2\theta + \alpha}{\theta} \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + (6 + \alpha) \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \frac{4 + \alpha}{2} \int_{\mathbb{R}^2} |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx$$

$$= (2 + \alpha) \gamma_{\mu}(a) + o_n(1).$$

Moreover, by $I_{\eta}(u_n) = \gamma_{\mu}(a) + o_n(1)$, we have

$$\frac{6+\alpha}{2} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u_{n})) F(u_{n}) dx + (6+\alpha) \gamma_{\mu}(a) + o_{n}(1)
= \eta \frac{6+\alpha}{\theta} \int_{\mathbb{R}^{2}} |\nabla u_{n}|^{\theta} dx + (6+\alpha) \int_{\mathbb{R}^{2}} |u_{n}|^{2} |\nabla u_{n}|^{2} dx + \frac{6+\alpha}{2} \int_{\mathbb{R}^{2}} |\nabla u_{n}|^{2} dx
\ge \eta \frac{2(\theta+\alpha)}{\theta} \int_{\mathbb{R}^{2}} |\nabla u_{n}|^{\theta} dx + (6+\alpha) \int_{\mathbb{R}^{2}} |u_{n}|^{2} |\nabla u_{n}|^{2} dx + \frac{4+\alpha}{2} \int_{\mathbb{R}^{2}} |\nabla u_{n}|^{2} dx
= \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u_{n})) f(u_{n}) u_{n} dx + (2+\alpha) \gamma_{\mu}(a) + o_{n}(1).$$

As $k > 3 + \frac{\alpha}{2}$, $2 < \theta < 3$ and (f_3) , we have that

$$\limsup_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx \le \frac{4}{k - (3 + \frac{\alpha}{2})} \gamma_\mu(a).$$

Lemma 3.5. The sequence $\{u_n\}$ satisfies

$$\limsup_{n \to +\infty} \left(\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \right) \le \theta(\frac{k-1-\frac{\alpha}{2}}{k-3-\frac{\alpha}{2}})\gamma_{\mu}(a).$$

Hence, there exists $\mu^* > 0$ such that

$$\limsup_{n\to +\infty} \left(\eta \int_{\mathbb{R}^2} |\nabla u_n|^\theta dx + \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \right) < \frac{(2+\alpha)\pi}{\gamma_0} - a^2 \quad \text{for} \quad \forall \mu \geq \mu^*.$$

Proof. Since $I_{\eta}(u_n) = \gamma_{\mu}(a) + o_n(1)$, we have

$$\lim_{n \to +\infty} \sup_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) F(u_n) dx + 2\gamma_{\mu}(a) + o_n(1)$$

$$= \eta \frac{2}{\theta} \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + 2 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx$$

$$> \frac{2}{\theta} \left(\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \right).$$

Therefore, by lemma 3.4, we have

$$\limsup_{n \to +\infty} \left(\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \right) \le \theta\left(\frac{k-1-\frac{\alpha}{2}}{k-3-\frac{\alpha}{2}}\right) \gamma_{\mu}(a).$$

By lemma 3.3, we can directly obtain the second inequality, the prove is complete.

Lemma 3.6. Fix $\mu \ge \mu^*$, where μ^* is given in Lemma 3.5. Then, $\{\lambda_n\}$ is a bounded sequence with

$$\limsup_{n \to +\infty} |\lambda_n| \le \left(\frac{8\theta(k - 1 - \frac{\alpha}{2}) + 2(2 + \alpha)}{a^2(k - 3 - \frac{\alpha}{2})} + \right) \gamma_{\mu}(a)$$

and

$$\liminf_{n \to +\infty} \lambda_n > \frac{2+\alpha}{2a^2} \liminf_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx.$$

Proof. From Lemma 2.6 and Lemma 3.5, we know that $\{u_n\}$ is bounded in O_r , and the boundedness of $\{u_n\}$ yields that $\{\lambda_n\}$ is bounded, indeed

$$\lambda_n |u_n|_2^2 = -\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n dx + o_n(1),$$

as for $|u_n|_2^2 = a^2$, we have

$$\lambda_n a^2 = -\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) f(u_n) u_n dx + o_n(1).$$
(3.10)

Hence,

$$|\lambda_n|a^2 \le \eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) f(u_n) u_n dx + o_n(1).$$

The limit (3.9) together with Lemma 3.4 and Lemma 3.5 ensure that the sequence $\left\{ \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n dx \right\}$ is bounded, because

$$\lim_{n \to +\infty} \sup \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u_{n})) f(u_{n}) u_{n} dx$$

$$< \lim_{n \to +\infty} \sup \left[4 \left(\eta \int_{\mathbb{R}^{2}} |\nabla u_{n}|^{\theta} dx + \int_{\mathbb{R}^{2}} |u_{n}|^{2} |\nabla u_{n}|^{2} dx + \int_{\mathbb{R}^{2}} |\nabla u_{n}|^{2} dx \right) + \frac{2+\alpha}{2} \int_{\mathbb{R}^{2}} (I_{\alpha} * F(u_{n})) F(u_{n}) dx \right]$$

$$< \frac{8\theta(k-1-\frac{\alpha}{2})+2(2+\alpha)}{k-3-\frac{\alpha}{2}} \gamma_{\mu}(a).$$

which concludes that $\{\lambda_n\}$ is a bounded sequence with

$$\limsup_{n \to +\infty} |\lambda_n| \le \left(\frac{8\theta(k-1-\frac{\alpha}{2}) + 2(2+\alpha)}{a^2(k-3-\frac{\alpha}{2})} + \right) \gamma_{\mu}(a)$$

In order to proof the second equality, the equality (3.10) together with the limit (3.9) lead to

$$\begin{split} \lambda_n a^2 &= -\eta \int_{\mathbb{R}^2} |\nabla u_n|^\theta dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n dx \\ &+ o_n(1) \\ &> -\eta \frac{2(\theta-1)}{\theta} \int_{\mathbb{R}^2} |\nabla u_n|^\theta dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n dx \\ &+ o_n(1) \\ &= \frac{2+\alpha}{2} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx + o_n(1), \end{split}$$

showing the desired result.

For fixed $\eta \in (0,1]$ and $\forall \mu \geq \mu^*$, from lemma 3.5, we can conclude that

$$\limsup_{n \to +\infty} |\nabla u_n|_2^2 < \frac{(2+\alpha)\pi}{\gamma_0} - a^2.$$

According to Corollary 2.3, we deduce that

$$\lim_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n dx = \int_{\mathbb{R}^2} (I_\alpha * F(u_\eta)) f(u_\eta) u_\eta dx$$

and

$$\lim_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx = \int_{\mathbb{R}^2} (I_\alpha * F(u_\eta)) F(u_\eta) dx,$$

where $u_n \to u_\eta$ in $H_r^{1,2}(\mathbb{R}^2)$. The last limit implies that $u_\eta \neq 0$, because otherwise, by Corollary 2.3 and the limit equation (3.9), we have

$$\lim_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx = \lim_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) f(u_n) u_n dx = 0,$$

and by lemma 3.6, we derive that

$$\liminf_{n \to +\infty} \lambda_n \ge \frac{2+\alpha}{2a^2} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx = 0.$$

Thus, we have $\lambda_n \geq 0$. Since $\{u_n\}$ is bounded in $H_r^{1,2}(\mathbb{R}^2)$. Corollary 2.3 together with (f_1) and (f_2) and the following equality

$$\lambda_n |u_n|_2^2 = -\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) f(u_n) u_n dx + o_n(1),$$

leads to

$$\lambda_n a^2 = -\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx - 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx - \int_{\mathbb{R}^2} |\nabla u_n|^2 dx + o_n(1).$$

From this, one has

$$0 \geq -\liminf_{n \to +\infty} \lambda_n a^2 = \limsup_{n \to +\infty} (-\lambda_n) a^2$$

$$= \limsup_{n \to +\infty} \left(\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \right)$$

$$\geq \liminf_{n \to +\infty} \left(\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \right)$$

$$\geq 0,$$

then, we obtain that

$$\eta \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx + 4 \int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx + \int_{\mathbb{R}^2} |\nabla u_n|^2 dx \to 0,$$

and this is impossible, because $\gamma_{\mu}(a) > 0$.

4 Critical points of $I_{\eta}|_{S(a)}$

The above analysis ensure that the weak limit u of $\{u_n\}$ is nontrivial. Moreover, the equality

$$\liminf_{n \to +\infty} \lambda_n \ge \frac{2+\alpha}{2a^2} \liminf_{n \to +\infty} \int_{\mathbb{R}^2} (I_\alpha * F(u_n)) F(u_n) dx$$

ensures that

$$\liminf_{n \to +\infty} \lambda_n = \frac{2+\alpha}{2a^2} \int_{\mathbb{R}^2} (I_\alpha * F(u_\eta)) F(u_\eta) dx > 0.$$

From this, going up to subsequence, still denoted by $\{\lambda_n\}$, we can assume that

$$\lambda_n \to \lambda_\eta > 0$$
, as $n \to +\infty$.

Since $\{u_n\}$ is bounded, we have

$$I'_{\eta}(u_n) + \lambda_{\eta} u_n \to 0 \quad in \ O^*.$$

Then, from [16, Lemma A.2], we have

$$I'_{\eta}(u_{\eta}) + \lambda_{\eta} u_{\eta} = 0, \tag{4.1}$$

testing (4.1) with $x \cdot \nabla u_{\eta}$ and u_{η} , we obtain $P_{\eta}(u_{\eta}) = 0$. It follows that

$$P_{\eta}(u_n) + \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) f(u_n) u_n dx - \frac{2+\alpha}{2} \int_{\mathbb{R}^2} (I_{\alpha} * F(u_n)) F(u_n) dx$$
$$\rightarrow P_{\eta}(u_{\eta}) + \int_{\mathbb{R}^2} (I_{\alpha} * F(u_{\eta})) f(u_{\eta}) u_{\eta} dx - \frac{2+\alpha}{2} \int_{\mathbb{R}^2} (I_{\alpha} * F(u_{\eta})) F(u_{\eta}) dx.$$

Then using the weak lower semicontinuous property, see [10, Lemma 4.3], there must be

$$\eta \frac{2(\theta - 1)}{\theta} \int_{\mathbb{R}^2} |\nabla u_n|^{\theta} dx \to \eta \frac{2(\theta - 1)}{\theta} \int_{\mathbb{R}^2} |\nabla u_\eta|^{\theta} dx, \tag{4.2}$$

$$\int_{\mathbb{R}^2} |u_n|^2 |\nabla u_n|^2 dx \to \int_{\mathbb{R}^2} |u_\eta|^2 |\nabla u_\eta|^2 dx, \tag{4.3}$$

$$\int_{\mathbb{R}^2} |\nabla u_n|^2 dx \to \int_{\mathbb{R}^2} |\nabla u_\eta|^2 dx. \tag{4.4}$$

That gives $I_{\eta}(u_{\eta}) = \lim_{n \to +\infty} I_{\eta}(u_n) = \gamma_{\mu}(a)$. Moreover, from (4.2)–(4.4), we obtain

$$I'_{n}(u_{n})[u_{n}] \to I'_{n}(u_{n})[u_{\eta}].$$
 (4.5)

Thus combining (4.5) with (4.1), there holds $\lambda_n |u_n|_2^2 \to \lambda_\eta |u_\eta|_2^2$. Since $\lambda_\eta > 0$, the last limit implies that $u_n \to u_\eta$ in O, implying that $|u_\eta|_2^2 = a^2$.

Based on the above preliminary works, we conclude that

Theorem 4.1. For any fixed $\eta \in (0,1]$, there exists a $u_{\eta} \in O_r \setminus \{0\}$ and a $\lambda_{\eta} \in \mathbb{R}$ such that

$$I'\eta(u_{\eta}) + \lambda_{\eta}u_{\eta} = 0,$$

 $I_{\eta}(u_{\eta}) = \gamma_{\mu}(a), \quad P_{\eta}(u_{\eta}) = 0,$
 $|u_{\eta}|_{2}^{2} = a^{2}.$

5 Proof of Theorem 1.1

By Theorem 4.1, we can take

$$\eta_n \to 0^+, \ I'_{\eta_n}(u_{\eta_n}) - \lambda_{\eta_n} u_{\eta_n} = 0, \text{ and } I_{\eta_n}(u_{\eta_n}) \to d^*(a) := \lim_{\eta_n \to 0^+} \gamma_{\mu}(a) \in (0, +\infty),$$

for $u_{\eta_n} \in S_r(a)$ with $|u_{\eta_n}|_2^2 = a^2$, then Lemma 3.6 implies that $\lambda_{\eta_n} < 0$. Now Theorem ?? gives that there exists $v \neq 0, v \in H^{1,2}_{rad}(\mathbb{R}^2) \cap L^{\infty}(\mathbb{R}^2)$, and $\lambda_0 \in \mathbb{R}$ such that

$$I'(v) - \lambda_0 v = 0$$
, $I(v) = d^*(a)$, and $|v|_2^2 = a^2$.

That is, v is a nontrivial radial solution of (1.4).

References

- [1] Agueh, M.: Sharp Gagliardo-Nirenberg inequalities via p-Laplacian type equations. NoDEA Non-linear Differential Equations Appl. 15, 457–472 (2008)
- [2] Alves, C.O., Figueiredo, G.M.: On multiplicity and concentration of positive solutions for a class of quasilinear problems with critical exponential growth in \mathbb{R}^N . J. Differential Equations. **246**, 1288–1311 (2009)
- [3] Alves, C.O., Ji, C., Miyagaki, O.H.: Normalized solutions for a Schrödinger equation with critical growth in \mathbb{R}^N . Calc. Var. Partial Differential Equations. **61**, 18 (2022)
- [4] Bartsch, T., Soave, N.: A natural constraint approach to normalized solutions of nonlinear Schrödinger equations and systems. J. Funct. Anal. 272, 4998–5037 (2017)
- [5] Berestycki, H., Lions, P.L.: Nonlinear scalar field equations. I. Existence of a ground state. *Arch. Rational Mech. Anal.* **82**, 313-345 (1983)
- [6] Berestycki, H., Lions, P.L.: Nonlinear scalar field equations. II. Existence of infinitely many solutions. Arch. Rational Mech. Anal. 82, 347–375 (1983)
- [7] Cao, D.M.: Nontrivial solution of semilinear elliptic equation with critical exponent in \mathbb{R}^2 Comm. Partial Differential Equations. 17, 407–435 (1992)
- [8] Cazenave, T., Lions, P.L.: Orbital stability of standing waves for some nonlinear Schrödinger equations Comm Math Phys. 85, 549–561 (1982)
- [9] Colin, M., Jeanjean, L.: Solutions for a quasilinear Schrödinger equation: a dual approach *Non-linear Anal.* **56**, 213–226 (2004)
- [10] Colin, M., Jeanjean, L., Squassina, M.: Stability and instability results for standing waves of quasi-linear Schrödinger equations. *Nonlinearity*. **23**, 1353–1385 (2010)
- [11] Dou, J.B., Huang, L., Zhong, X.X.: Normalized Solutions to N-Laplacian Equations in \mathbb{R}^N with Exponential Critical Growth. Journal of Geometric Analysis. **34**, 317 (2024)
- [12] Ghoussoub, N.: Duality and perturbation methods in critical point theory. Cambridge Tracts in Mathematics. 107, (1993) (ISBN: 0-521-44025-4)

- [13] Hirata, J., Tanaka, K.: Nonlinear scalar field equations with L^2 constraint mountain pass and symmetric mountain pass approaches. Adv. Nonlinear Stud. 19, 263–290 (2019)
- [14] Ikoma, N., Tanaka, K.: A note on deformation argument for L^2 normalized solutions of nonlinear Schrödinger equations and systems. Adv. Differential Equations. 24, 609–646 (2019)
- [15] Jeanjean, L.: Existence of solutions with prescribed norm for semilinear elliptic equations. Nonlinear Anal. 28, 1633–1659 (1997)
- [16] Li, H.W., Zou, W.M.: Quasilinear Schrödinger equations: ground state and infinitely many normalized solutions. *Pacific J. Math.* **322** , 99–138 (2023)
- [17] Litvak, A.G., Sergeev, A.M.: One dimensional collapse of plasma waves. *JETP*, *Letters.* **27**, 517–520 (1978)
- [18] Liu, X.Q., Liu, J.Q., Wang, Z.Q.: Ground states for quasilinear Schrödinger equations with critical growth. Calc. Var. Partial Differential Equations. 46, 641–669 (2013)
- [19] Liu, J.Q., Wang, Z.Q.: Multiple solutions for quasilinear elliptic equations with a finite potential well. J. Differential Equations. 257, 2874–2899 (2014)
- [20] Liu, J.Q., Wang, Y.Q., Wang, Z.Q.: Soliton solutions for quasilinear Schrödinger equations. II. J. Differential Equations. 187, 473–493 (2003)
- [21] Liu, X.Q., Liu, J.Q., Wang, Z.Q.: Quasilinear elliptic equations via perturbation method. *Proc. Amer. Math. Soc.* **141**, 253–263 (2013)
- [22] Liu, X.Q., Liu, J.Q., Wang, Z.Q.: Quasilinear elliptic equations with critical growth via perturbation method *J. Differential Equations.* **254**, 102–124 (2013)
- [23] Makhankov, V.G., Fedyanin, V.K.: Nonlinear effects in quasi-one-dimensional models and condensed matter theory Phys. Rep. 104, 1–86 (1984)
- [24] Moser, J.: A sharp form of an inequality by N. Trudinger. *Indiana Univ. Math. J.* **20**, 1077–1092 (1970)
- [25] Palais, R.S.: The principle of symmetric criticality. Comm. Math. Phys. 69, 19–30 (1979)
- [26] Poppenberg, M., Schmitt, K., Wang, Z.Q.: On the existence of soliton solutions to quasilinear Schrödinger equations. Calc. Var. Partial Differential Equations. 14, 329–344 (2002)
- [27] Porkolab, M., Goldman, M.V.: Upper-hybrid solitons and oscillating-two-stream instabilities. *Phys. Fluids.* **19**, 872–881 (1976)
- [28] Shibata, M.: Stable standing waves of nonlinear Schrödinger equations with a general nonlinear term. *Manuscripta Math.* **143**, 221–237 (2014)
- [29] Soave, N.: Normalized ground states for the NLS equation with combined nonlinearities. *J. Dif*ferential Equations. **269**, 6941–6987 (2020)
- [30] Soave, N.: Normalized ground states for the NLS equation with combined nonlinearities: the Sobolev critical case. J. Funct. Anal. 279, 108610 (2020)

- [31] Stuart, A.: Bifurcation from the continuous spectrum in the L^2 -theory of elliptic equations on \mathbb{R}^n . Recent methods in nonlinear analysis and applications (Naples, 1980). 231–300 (1981)
- [32] Stuart, C.A.: Bifurcation in $L^p(\mathbf{R}^N)$ for a semilinear elliptic equation. *Proc. London Math. Soc.* (3). **57**, 511–541 (1988)
- [33] Stuart, C.A.: Bifurcation from the essential spectrum for some noncompact nonlinearities. *Math. Methods Appl. Sci.* **11**, 525–542 (1989)
- [34] Trudinger, N.S.: On imbeddings into Orlicz spaces and some applications. J. Math. Mech. 17, 473–483 (1967)
- [35] Willem, M.: Minimax Theorems. Birkhäuser Boston, Inc., Boston, MA. (1996) (ISBN: 0-8176-3913-6)